

Non-Perturbative Effects & Heterotic String Vacua

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Jacob M. Leedom
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CLUSTER OF EXCELLENCE
QUANTUM UNIVERSE

Overview

- > Work with toroidal orbifold compactifications of the Heterotic string, e.g. $X_6 = T^6/\mathbb{Z}_N$
- > Prove no-go theorem for dS minima for Kähler modulus + dilaton sector including non-perturbative effects in superpotential
- > Attempt to extend no-go results to stronger non-perturbative effects in g_s & indicate potential loophole

No-Go Results

- > Classical SUGRA [Maldacena,Nunez,'00]: no dS
- > Subleading α' corrections [Green,Martinec,Quigley,Sethi,'11]: AdS, but no dS
- > Infinite tower of α' [Gautason,Junghans,Zagermann,'12]: no AdS or dS
- > All worldsheet effects [Kutasov,Maxfield,Melnikov,Sethi,'15]: no dS
- > Gaugino Condensation [Quigley,'15]: no AdS or dS
- > Gaugino Condensation, threshold effects, instantons* [Gonzalo, Ibanez,Uranga,'18]: AdS, numerically no dS

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$$\delta\mathcal{L} \sim e^{-S} = e^{-1/g_s^2}$$

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$$\delta\mathcal{L} \sim e^{-S} = e^{-1/g_s^2}$$

What about non-perturbative effects of strength e^{-1/g_s} ?

Two-Moduli Model

> Kähler Modulus T and Dilaton S

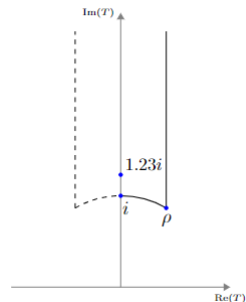
- T has an $SL(2, \mathbb{Z})$ duality:

$$T \rightarrow \gamma \cdot T = \frac{aT + b}{cT + d}$$

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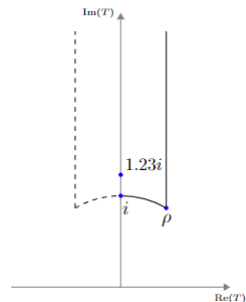
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- Kähler potential

$$\mathcal{K} = -\ln(S + \bar{S}) - 3\ln(i(T - \bar{T}))$$

- Superpotential

$$W(S, T) = \frac{\Omega(S)H(T)}{\eta^6(T)}$$



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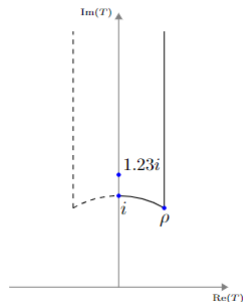
$$W(S, T) = \frac{\Omega(S)H(T)}{\eta^6(T)}$$

$$\Omega(S) = h + \sum_a \Lambda_a^3 e^{-k_a S/b_a}$$

$$H(T) = \left(\frac{G_4(T)}{\eta^8(T)} \right)^n \left(\frac{G_6(T)}{\eta^{12}(T)} \right)^m \mathcal{P}(j(T))$$

[Rademacher,Zuckerman,'38]

[Discontinuous Groups and Automorphic Functions - Lehner]



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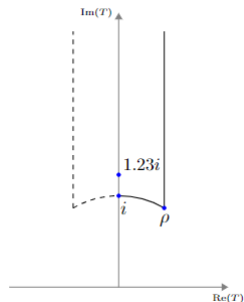
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- Scalar Potential

$$V(S, T) = e^{\mathcal{K}} (K^{S\bar{S}} F_S \bar{F}_{\bar{S}} + K^{T\bar{T}} F_T \bar{F}_{\bar{T}} - 3\bar{W}W)$$

$$= e^{k(S, \bar{S})} Z(T, \bar{T}) |\Omega(S)|^2 \left\{ |H(T)|^2 (A(S, \bar{S}) - 3) + \hat{V}(T, \bar{T}) \right\}$$



$$A(S, \bar{S}) = \frac{K^{S\bar{S}} F_S \bar{F}_{\bar{S}}}{|W|^2} = \frac{K^{S\bar{S}} |\Omega_S + K_s \Omega|^2}{|\Omega|^2}$$

$$Z(T, \bar{T}) = \frac{1}{i(T - \bar{T})^3 |\eta(T)|^{12}}$$

$$\hat{V}(T, \bar{T}) = \frac{-(T + \bar{T})^2}{3} \left| H'(T) - \frac{3i}{2\pi} H(T) \hat{G}_2(T, \bar{T}) \right|^2$$

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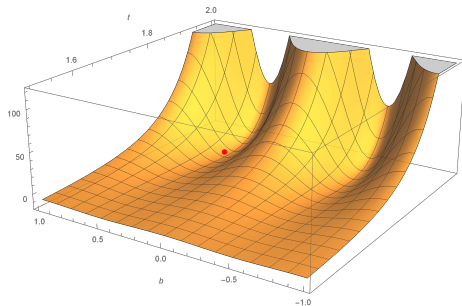
$$\mathcal{K} = -k(S, \bar{S}) - 3 \ln(i(T - \bar{T}))$$

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A dS No-Go Result

- > Assume typical Heterotic dilaton stabilization:

$$\langle F_S \rangle = 0 \rightarrow A(S, \bar{S}) = 0$$

- > Conditions for dS minimum at $T = T_0$:

$$V_0 = \Lambda^4 > 0$$

$$\partial_T V_0 = 0$$

$$\text{Hes}(V)_0 \ \& \ \partial_t^2 V_0 > 0$$

- > Solve for H'_0 & H''_0 , plug into 2nd derivatives:

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Cannot both be positive



dS minima not possible

Beyond the No-Go

- > $V(T, S)$ is a modular function in T , so $\partial_T V$ is a weight 2 modular form and vanishes at $T = i, \rho$.
- > Self dual points always extremum - when are they minima?

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- > Set $n = 0$ in parametrization of $H(T)$. At $T = \rho$:

$$V(S, \bar{S}, \rho, \rho^*) = \frac{4^{4m+6} \pi^{12m+12}}{1225^m \times 27^{m+1} \Gamma^{18}(1/3)} |\Omega(S)|^2 |\mathcal{P}(0)|^2 e^{k(S, \bar{S})} (A(S, \bar{S}) - 3)$$
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One method: Shenker-like Non-Perturbative Effects

Non-perturbative Contributions: Shenker-like Effects

- > All closed string theories have non-perturbative contributions of strength e^{-1/g_s}
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$$\begin{array}{lll} \text{Type I Worksheet Instantons} & \delta\mathcal{L}_I \sim e^{-A^I/\alpha'} & \leftrightarrow \delta\mathcal{L}_H \sim e^{-\frac{A^H}{\alpha'\lambda}} \\ \text{Type IIA Worldline Instantons} & \delta\mathcal{L}_{IIA} \sim \sum_m e^{-mR^{IIA}} & \leftrightarrow \delta\mathcal{L}_H \sim \sum_m e^{-m/g_s} \end{array}$$

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- > Also in M-theory 1-loop amplitudes [\[Green & Rudra, '16\]](#)

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- > Also in M-theory 1-loop amplitudes [Green & Rudra, '16]

$$S^{HO} \supset \frac{g_s^{-1/2}}{2^9 (2\pi)^7 4! \ell_H^2} \int_{\mathcal{M}_{10}} \sqrt{-G} t_8 t_8 R^4 E_{3/2}(g_s^{-1})$$
$$E_{\frac{3}{2}}(g_s) = 2\zeta(3)g_s^{-3/2} + 2\zeta(2)g_s^{1/2} + \sum_{n \in \mathbb{Z}^+} 8\pi\sigma_{-1}(|n|)e^{-\frac{2\pi|n|}{g_s}}(1 + \mathcal{O}(g_s))$$

A Loophole Example

> Linear Multiplet Formalism: $L \supset \{\ell, \psi, B_2\}$

$$\langle \ell \rangle = \frac{g_s^2}{2}$$

$$\mathcal{L}_{KE} = \int d^4\theta E \left(-2 + f(L) \right) \quad k(L) = \ln(L) + g(L)$$

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> Parametrize Shenker-like effects [Gaillard & Nelson,'07]+:

$$f(\ell) = \sum_{n=0} A_n \ell^{-q_n} e^{-B/\sqrt{\ell}} \quad L \frac{df}{dL} = -L \frac{dg}{dL} + f$$

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> Scalar potential for single gaugino condensate:

$$V(\ell) = \frac{\mathcal{T}}{\ell} \left[(1 + \ell g')(1 + b\ell)^2 - 3b^2 \ell^2 \right] e^{g - (f+1)/b\ell}$$

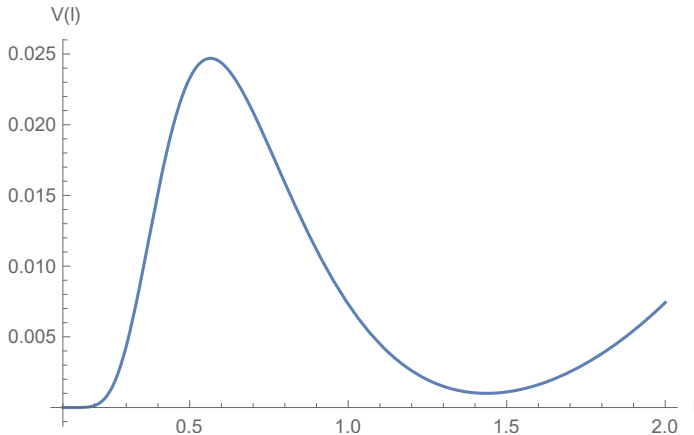
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$$\left\langle \frac{\ell}{1 + f(\ell)} \right\rangle = \frac{g_s^2}{2}$$

$$A_0 = 27$$

$$B = \pi$$

$$b_{E_8} = \frac{30}{8\pi^2}$$

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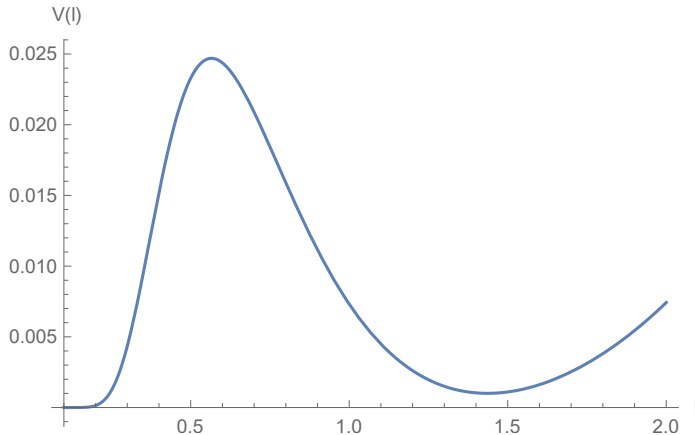
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$$\langle e^{-B/\sqrt{\ell}} \rangle \sim 0.07$$

$$\langle e^{-2B/\sqrt{\ell}} \rangle \sim 0.005$$

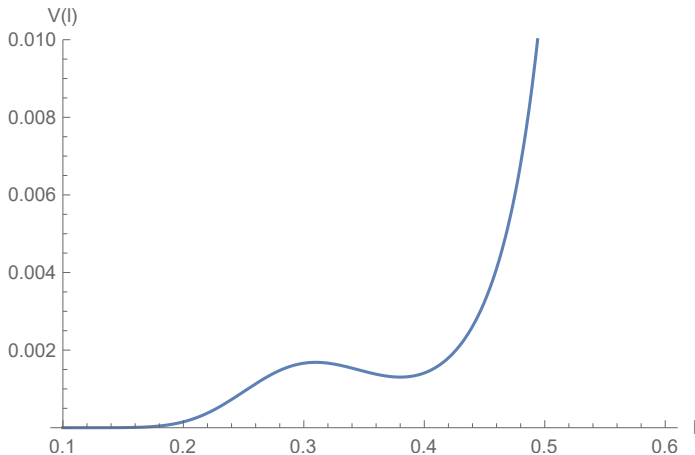
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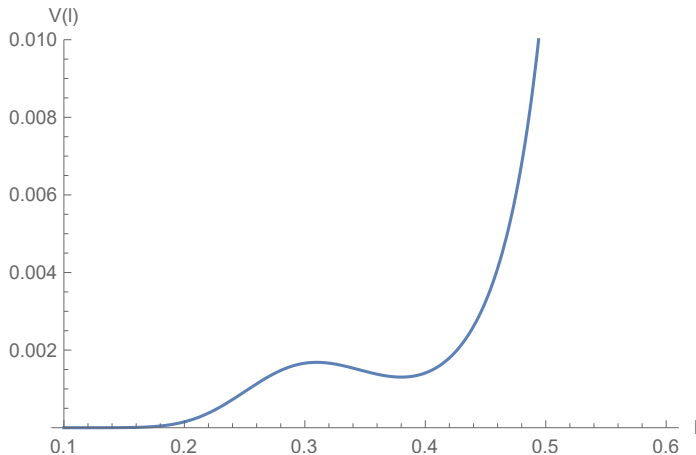
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$$\langle e^{-B/\sqrt{\ell}} \rangle \sim 3.7 \times 10^{-5}$$

$$\langle e^{-2B/\sqrt{\ell}} \rangle \sim 1.4 \times 10^{-9}$$

Conclusions

- > Provided a no-go theorem $\Rightarrow$ extends connection between modular-invariance and swampland conjectures
- > Illustrated a loophole in dS no-go arguments using Shenker-like e^{-1/g_s} effects in the Kähler potential
- > Many potential future directions - most obvious is a better understanding of Shenker-like effects to close loophole or realize meta-stable dS in Heterotic string constructions
- > Silverstein-like Effects - $\delta\mathcal{L} \sim e^{-1/g_s^{1/2}}$

Thank you!

Contact

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